

# Exit time asymptotics on non-commutative 2-torus.

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The purpose of this talk is to establish an analogue of exit time asymptotics of Brownian motion on manifolds, in the set-up of non-commutative 2-torus. Using these asymptotics, we will try to formulate definitions of certain geometric invariants e.g. intrinsic dimension, mean curvature etc for the non-commutative 2-torus.

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- 1** Interplay between Geometry and Probability:
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- 3** A case study:Exit time asymptotics on the non-commutative 2-torus

# Exit time asymptotics of Brownian motion on manifolds:

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We begin with the following well-known proposition:

Pinsky, 1994

Consider a hypersurface  $M \subseteq \mathbb{R}^d$  with the Brownian motion process  $X_t^m$  starting at  $m$ . Let  $T_\varepsilon = \inf\{t > 0 : \|X_t^m - m\| = \varepsilon\}$  be the exit time of the motion from an extrinsic ball of radius  $\varepsilon$  around  $m$ . Then we have

$$\mathbb{E}_m(T_\varepsilon) = \varepsilon^2/2(d-1) + \varepsilon^4 H^2/8(d+1) + O(\varepsilon^5),$$

where  $H$  is the mean curvature of  $M$ .

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## Gray, 1973

Let  $V_m(\epsilon)$  denote the volume of a ball of radius  $\epsilon$  around  $m \in M$ . Let  $n$  be the intrinsic dimension of the manifold. Then we have

$$V_m(\epsilon) = \frac{\alpha_n \epsilon^n}{n} \left( 1 - K_1 \epsilon^2 + K_2 \epsilon^4 + O(\epsilon^6) \right)_m,$$

where  $\alpha_n := 2\Gamma(\frac{1}{2})^n \Gamma(\frac{n}{2})^{-1}$  and  $K_1, K_2$  are constants depending on the manifold.

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The intrinsic dimension  $n$  of the hypersurface  $M$  is the unique integer  $n$

$$\text{satisfying } \lim_{\epsilon \rightarrow 0} \frac{\mathbb{E}(\tau_\epsilon)}{V_\epsilon^m} = \begin{cases} \infty & \text{if } m \text{ is less than } n; \\ \neq 0 & \text{if } m \neq n; \\ = 0 & \text{if } m > n. \end{cases}$$

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Observe that  $\frac{V(\epsilon)^{\frac{2}{n}}}{\epsilon^2} \rightarrow \left(\frac{\alpha_n}{n}\right)^{\frac{2}{n}}$  and  $\frac{V(\epsilon)^{\frac{4}{n}}}{\epsilon^4} \rightarrow \left(\frac{\alpha_n}{n}\right)^{\frac{4}{n}}$  as  $\epsilon \rightarrow 0^+$ .

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In view of this, the asymptotic expression appearing in Pinsky's result can be recast as

$$\mathbb{E}(\tau_\epsilon) = \frac{1}{2(d-1)} \left(\frac{V(\epsilon)n}{\alpha_n}\right)^{\frac{2}{n}} + \frac{H^2}{8(d+1)} \left(\frac{V(\epsilon)n}{\alpha_n}\right)^{\frac{4}{n}} + O(V(\epsilon)^{\frac{5}{n}}).$$

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$$d = \frac{1}{2} \left(1 + \lim_{\epsilon \rightarrow 0} \frac{1}{\mathbb{E}(\tau_\epsilon)} \left(\frac{nV(\epsilon)}{\alpha_n}\right)^{\frac{2}{n}}\right), \quad (1)$$

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$$H^2 = 8(d+1) \left(\frac{\alpha_n}{n}\right)^{\frac{4}{n}} \lim_{\epsilon \rightarrow 0} \frac{\mathbb{E}(\tau_\epsilon) - \frac{1}{2(d-1)} \left(\frac{nV(\epsilon)}{\alpha_n}\right)^{\frac{2}{n}}}{V(\epsilon)^{\frac{4}{n}}}. \quad (2)$$

# Formulation of quantum exit time.

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$$\chi_{\{\tau_{B_r^x} > t\}} = \bigwedge_{s \leq t} \left( \chi_{\{W_s^x \in B_r^x\}} \right),$$

where  $\bigwedge$  denotes infimum.

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Note that

$$j_t : L^\infty(U_x) \rightarrow L^\infty(U_x) \otimes B(\Gamma(L^2(\mathbb{R}_+, \mathbb{C}^n))),$$

since by the Wiener- Itô isomorphism,  $L^2(\mathbb{P}) \cong \Gamma(L^2(\mathbb{R}_+, \mathbb{C}^n))$ , where  $\mathbb{P}$  is the  $n$  dimensional Wiener measure.

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So one may write

$$\chi_{\{\tau_{B_r^x} > t\}}(\cdot) = \bigwedge_{s \leq t} j_s(\chi_{B_r^x})(x, \cdot) = \bigwedge_{s \leq t} ((ev_x \otimes id) \circ j_s(\chi_{B_r^x}))(\cdot).$$

Thus we may view  $\tau_{B_r^x}$  as a spectral family in  $L^\infty(U_x) \otimes B(\Gamma(L^2(\mathbb{R}_+, \mathbb{C}^n)))$  by the prescription:

$$\tau_{B_r^x}([0, t]) = \mathbf{1} - \wedge_{s \leq t} (j_s(\chi_{B_r^x})) .$$

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Moreover, we have:

$$\mathbb{E}(\tau_{B_r^x}) = \int_0^\infty \mathbb{P}(\tau_{B_r^x} > t) dt = \int_0^\infty \langle e(0), \{ (ev_x \otimes \mathbf{1}) (\wedge_{s \leq t} j_s(\chi_{B_r^x})) \} e(0) \rangle dt .$$

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Alternatively:

Choose a sequence  $(x_n)_n \in M$  and positive numbers  $\epsilon_n$  such that  $x_n \rightarrow x$  and  $\epsilon_n \rightarrow 0$ . Now for large  $n$ ,  $\chi_{\{W_s^{x_n} \in B_{\epsilon_n}^{x_n}\}}(\cdot) \stackrel{\mathcal{L}}{=} \chi_{\{W_s^x \in B_{\epsilon_n}^x\}}(\cdot)$  for each  $s \geq 0$ . Thus,

$$\mathbb{E}(\tau_{B_{\epsilon_n}^{x_n}}) = \int_0^\infty \langle e(0), \{(ev_{x_n} \otimes id) (\wedge_{s \leq t} j_s(\chi_{B_{\epsilon_n}^{x_n}}))\} e(0) \rangle dt = \mathbb{E}(\tau_{B_{\epsilon_n}^x}),$$

i.e. the asymptotic behaviour of  $\mathbb{E}(\tau_{B_{\epsilon_n}^{x_n}})$  and  $\mathbb{E}(\tau_{B_{\epsilon_n}^x})$  will be the same.

Note that the points of  $M$  are in 1 – 1 correspondence with the pure states of  $L^\infty(M)$  and  $\{P_n = \chi_{B_{\epsilon_n}^{x_n}}\}_n$  is a family of projections on  $L^\infty(M)$ , so that we have:

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We now move into non-commutative setup.

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### Barnett, Wilde, 1991

Let  $(\mathfrak{A}_t)_{t \geq 0}$  be an increasing family of von-Neumann algebras (called a filtration). A quantum random time or stop time adapted to the filtration  $(\mathfrak{A}_t)_{t \geq 0}$  is an increasing family of projections  $(E_t)_{t \geq 0}$ ,  $E_0 = I$  such that  $E_t$  is a projection in  $\mathfrak{A}_t$  and  $E_s \leq E_t$  whenever  $0 \leq s \leq t < +\infty$ .

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Observe that by our definition,  $\tau_{B_f^x}([0, t])$  is adapted to the filtration

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$\mathfrak{A}_t := L^\infty(U_x) \otimes B(\Gamma_{[t]})$  ( $\Gamma_{[t]} := \Gamma(L^2([0, t], \mathbb{C}^n))$ ), for

$\tau_{B_f^x}([0, t]) \in \mathfrak{A}_t \otimes 1_{\Gamma_{[t]}}$ .

Suppose that we are given an E-H flow  $j_t : \mathcal{A} \rightarrow \mathcal{A}'' \otimes B(\Gamma(L^2(\mathbb{R}_+, k_0)))$ , where  $\mathcal{A}$  is a  $C^*$  or von-Neumann algebra. For a projection  $P \in \mathcal{A}$ , the family  $\{\mathbf{1} - \wedge_{s \leq t} (j_s(P))\}_{t \geq 0}$  defines a quantum random time adapted to the filtration  $(\overline{\mathcal{A}'' \otimes B(\Gamma_t)})_{t \geq 0}$ .

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## Definition

*We refer to the quantum random time  $\{1 - \bigwedge_{s \leq t} j_s(P)\}_{t \geq 0}$  as the 'exit time from the projection  $P$ .*

Let  $\tau$  be a state (to be thought of as non-commutative volume form on a  $C^*$  or von Neumann algebra), and assume that we are given a family  $\{P_n\}_{n \geq 1}$  of projections in  $\mathcal{A}$ , and a family  $\{\omega_n\}_{n \geq 1}$  of pure states of  $\mathcal{A}$  such that

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## Definition

Let  $\gamma_n := \int_0^\infty dt \langle e(0), (\omega_n \otimes id) \circ \bigwedge_{s \leq t} j_s(P_n) e(0) \rangle$ . We say that there is an exit time asymptotic for the family  $\{P_n; \omega_n\}$  of intrinsic dimension  $n_0$  if

$$\lim_{n \rightarrow \infty} \frac{\gamma_n}{v_n^m} = \begin{cases} \infty & \text{if } m \text{ is just less than } n_0 \\ \neq 0 & \text{if } m \neq n \\ = 0 & \text{if } m > n \end{cases}$$

and

$$\gamma_n = c_1 v_n^{\frac{2}{n_0}} + c_2 v_n^{\frac{4}{n_0}} + \cdots + c_k v_n^{\frac{2k}{n_0}} + O(v_n^{\frac{2k+1}{n_0}}) \text{ as } n \rightarrow \infty. \quad (3)$$

It is not at all clear whether such an asymptotic exists in general, and even if it exists, whether it is independent of the choice of the family  $\{P_n; \omega_n\}$ . If it is the case, one may legitimately think of  $c_1, c_2, \dots, c_k \dots$  as geometric invariants and imitating the classical formulae as discussed before, the extrinsic dimension  $d$  and the mean curvature  $H$  of the non-commutative manifold may be defined to be

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$$d := \frac{1}{2c_1} \left( \frac{n_0}{\alpha_{n_0}} \right)^{\frac{2}{n_0}} + 1, \quad (4)$$

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$$d := \frac{1}{2c_1} \left( \frac{n_0}{\alpha_{n_0}} \right)^{\frac{2}{n_0}} + 1, \quad (4)$$

$$H^2 := 8(d+1)c_2 \left( \frac{\alpha_{n_0}}{n_0} \right)^{\frac{4}{n_0}}. \quad (5)$$

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Fix an irrational number  $\theta \in [0, 1]$ .

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Fix an irrational number  $\theta \in [0, 1]$ .

## Definition

*The non-commutative 2-torus  $C^*(\mathbb{T}_\theta^2)$  is the universal  $C^*$ -algebra generated by a pair of unitaries  $U, V$  which satisfy:*

$$UV = e^{2\pi i\theta} VU.$$

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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It can also be viewed as the “Rieffel deformation” of the commutative  $C^*$ -algebra  $C(\mathbb{T}^2)$ .

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

A class of projections on  $C^*(\mathbb{T}_\theta^2)$ , as given by Rieffel, is:

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Choose an  $\epsilon \ll \theta$  and let  $P = f_{-1}(U)V^{-1} + f_0(U) + f_1(U)V$ , where  $f_1, f_0 \in C(\mathbb{T}^2)$ ,  $f_{-1}(t) := \overline{f_1(t + \theta)}$ ,

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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$$f_0(t) = \begin{cases} \epsilon^{-1}t & \text{if } 0 \leq t \leq \epsilon \\ 1 & \text{if } \epsilon \leq t \leq \theta \\ \epsilon^{-1}(\theta + \epsilon - t) & \text{if } \theta \leq t \leq \theta + \epsilon \\ 0 & \text{if } \theta + \epsilon \leq t \leq 1 \end{cases}$$

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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$$f_1(t) = \begin{cases} \sqrt{f_0(t) - f_0(t)^2} & \text{if } \theta \leq t \leq \theta + \epsilon \\ 0 & \text{if otherwise.} \end{cases}$$

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

- Let  $tr$  be the canonical trace in  $C^*(\mathbb{T}_\theta^2)$ , given by  $tr(\sum_{m,n} a_{mn} U^m V^n) = a_{00}$ . This trace will be taken as an analogue of the volume form in  $C^*(\mathbb{T}^2)$ .

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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- Throughout the section, we will assume  $C^*(\mathbb{T}_\theta^2) \subseteq B(L^2(tr))$ , and let  $W^*(\mathbb{T}_\theta^2) := (C^*(\mathbb{T}_\theta^2))''$ .

# Exit time asymptotics on the non-commutative 2-torus

B. Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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- For  $(x, y) \in \mathbb{T}^2$ , let  $\alpha_{(x,y)}$  denote the canonical action of  $\mathbb{T}^2$  on  $C^*(\mathbb{T}_\theta^2)$  given by  $\alpha_{(x,y)}(\sum_{m,n} a_{mn} U^m V^n) = \sum_{m,n} x^m y^n a_{mn} U^m V^n$ . Note that the automorphism  $\alpha$  is  $tr$ -preserving. Hence it extends to a unitary operator on  $L^2(tr)$ , say  $u_{(x,y)}$ , and  $\alpha = ad u$ , which implies that  $\alpha$  is normal.

# Exit time asymptotics on the non-commutative 2-torus

B. Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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- On  $C^*(\mathbb{T}_\theta^2)$ , there are two conditional expectations denoted by  $\phi_1, \phi_2$ , which are defined as:

$$\phi_1(A) := \int_0^1 \alpha_{(1, e^{2\pi i t})}(A) dt, \quad \phi_2(A) := \int_0^1 \alpha_{(e^{2\pi i t}, 1)}(A) dt.$$

From the normality of  $\alpha$ , it follows easily that  $\phi_1, \phi_2$  are normal maps.

# Exit time asymptotics on the non-commutative 2-torus

B. Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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From the normality of  $\alpha$ , it follows easily that  $\phi_1, \phi_2$  are normal maps.

- For a projection  $P$ , let  $A_{(s,t)}(P) := \alpha_{e^{2\pi i s}, e^{2\pi i t}}(P)$ .

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

## Theorem

Let  $P = f_{-1}(U)V^{-1} + f_0(U) + f_1(U)V$  be a projection such that  $f_0, f_1$  satisfy the conditions described before. Consider the projections  $A_{s,t}(P)$ ,  $A_{s',t'}(P)$  such that  $|s - s'| < \frac{\epsilon}{4}$ . Then

$$(A_{s,t}(P)) \wedge (A_{s',t'}(P)) = \chi_S(U),$$

for the set  $S = X_1 \cap X_2 \cap X_3 \cap X_4$ , where

$X_1 = \tau_{-s}(\{x | f_1(x) = 0\})$ ,  $X_2 := \tau_{-s'}(\{x | f_1(x) = 0\})$ ,

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It is worthwhile to note that the conclusion of the above theorem holds if we replace  $U$  by  $U^k$ ,  $V$  by  $V^k$ , and  $\theta$  by  $\{k\theta\}$  ( $\{\cdot\}$  denoting the fractional part).

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Let  $P_n = f_{-1}^{(k_n)}(U^{k_n}) + f_0^{(k_n)}(U^{k_n}) + f_1^{(k_n)}(U^{k_n})U^{k_n}$ , be projections such that  $\{k_n\theta\} \rightarrow 0$ . Put  $\epsilon := \frac{\{k_n\theta\}}{2}$ .

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Consider a standard Brownian motion in  $\mathbb{R}^2$ , given by  $(W_t^{(1)}, W_t^{(2)})$ .

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Consider a standard Brownian motion in  $\mathbb{R}^2$ , given by  $(W_t^{(1)}, W_t^{(2)})$ .

Define  $j_t : W^*(\mathbb{T}_\theta^2) \rightarrow W^*(\mathbb{T}_\theta^2) \otimes B(\Gamma(L^2(\mathbb{R}_+, \mathbb{C}^2)))$  by  
$$j_t(\cdot) := \alpha_{(e^{2\pi i W_t^{(1)}} , e^{2\pi i W_t^{(2)}})}(\cdot).$$

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Note that  $j_t$  defined above is the standard Brownian motion on  $C^*(\mathbb{T}_\theta^2)$ .

# Exit time asymptotics for non-commutative 2-torus

We have:

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

# Exit time asymptotics for non-commutative 2-torus

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B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

# Exit time asymptotics for non-commutative 2-torus

We have:

## Theorem

*Almost surely,  $\bigwedge_{s \leq t} (j_s(P_n)(\omega)) \in W^*(U)$ , for all  $n$ , i.e.*

$$\bigwedge_{s \leq t} (j_s(P_n)) \in W^*(U) \otimes B(\Gamma(L^2(\mathbb{R}_+, \mathbb{C}^2))),$$

*for each  $n$ .*

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
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for each  $n$ .

## Outline of the proof:

In the strong operator topology,

$$\bigwedge_{0 \leq s \leq t} (j_s(P_n)) = \lim_{m \rightarrow \infty} \bigwedge_i \{j_{\frac{it}{2^m}}(P_n) \wedge j_{\frac{(i+1)t}{2^m}}(P_n)\}. \quad (6)$$

Now almost surely a Brownian path restricted to  $[0, t]$  is uniformly continuous, so that for sufficiently large  $m$ , and for almost all  $\omega$ ,  $|W_{\frac{it}{2^m}}^{(1)} - W_{\frac{(i+1)t}{2^m}}^{(1)}|$  can be made small, uniformly for all  $i$  such that  $i = 0, 1, \dots, 2^m$ . So  $\bigwedge_i \{j_{\frac{it}{2^m}}(P_n) \wedge j_{\frac{(i+1)t}{2^m}}(P_n)\} \in W^*(U)$  by Theorem 3.2. It can be shown that the set of projections of this type is closed in the WOT-topology. Hence proved. ↻

B.Das

Interplay between Geometry and Probability:

Exit time asymptotics of Brownian motion on manifolds.

Formulation of quantum exit time.

A case study: Exit time asymptotics on the non-commutative 2-torus

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Note that  $W^*(U)$  is isomorphic with  $L^\infty(\mathbb{T})$ .

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Note that  $W^*(U)$  is isomorphic with  $L^\infty(\mathbb{T})$ .

Consider the pure states  $\{ev_z \circ E_1, ev_x \circ E_2 | x, z \in \mathbb{T}\}$  on  $W^*(\mathbb{T}_\theta^2)$ , which are also normal. Let  $z_n = e^{2\pi i \frac{3\{kn\theta\}}{4}}$ . Consider the sequence of pure states  $\phi_{z_n} := ev_{z_n} \circ E_1$ .

# Exit time asymptotics for non-commutative 2-torus

B. Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Note that  $W^*(U)$  is isomorphic with  $L^\infty(\mathbb{T})$ .

Consider the pure states  $\{ev_z \circ E_1, ev_x \circ E_2 | x, z \in \mathbb{T}\}$  on  $W^*(\mathbb{T}_\theta^2)$ , which are also normal. Let  $z_n = e^{2\pi i \frac{3\{kn\theta\}}{4}}$ . Consider the sequence of pure states  $\phi_{z_n} := ev_{z_n} \circ E_1$ .

Consider

$$\langle e(0), (\phi_{z_n} \otimes \mathbf{1}) \circ \bigwedge_{0 \leq s \leq t} (j_s(P_n))e(0) \rangle.$$

# Exit time asymptotics for non-commutative 2-torus

B. Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
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A direct computation shows that this is equal to

$$\mathbb{P}\{e^{2\pi i W_s^{(1)}} \in \mathcal{B}, 0 \leq s \leq t\} = \mathbb{P}\{\tau_{[-\frac{\{k_n\theta\}}{4}, \frac{\{k_n\theta\}}{4}]} > t\},$$

where  $\mathcal{B} := \{e^{2\pi i x} : x \in [-\frac{\{k_n\theta\}}{4}, \frac{\{k_n\theta\}}{4}]\}$ .

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B. Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Consider

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where  $\mathcal{B} := \{e^{2\pi i x} : x \in [-\frac{\{k_n\theta\}}{4}, \frac{\{k_n\theta\}}{4}]\}$ .

So we have a family of  $(\tau_n)_n$  random times defined by

$$\tau_n([t, +\infty)) = \bigwedge_{0 \leq s \leq t} (j_s(P_n));$$

so that  $\int_0^t \langle e(0), (\phi_{z_n} \otimes \mathbf{1}) \circ \bigwedge_{0 \leq s \leq t} (j_s(P_n))e(0) \rangle dt$  can be taken as the expectation of the random time  $\tau_n$ .

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Note that here the analogue for balls of decreasing volume is  $(P_n)_n$ , such that  $tr(P_n) = \{k_n\theta\} \rightarrow 0$ ,  $tr$  being the canonical trace in  $W^*(\mathbb{T}_\theta^2)$ .

# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Now, by the Pinsky's result, we have

$$\begin{aligned} & \int_0^t \langle e(0), (\phi_{z_n} \otimes 1) \circ \bigwedge_{0 \leq s \leq t} (j_s(P_n)) e(0) \rangle dt \\ &= \mathbb{E}(\tau_{[-\frac{\{k_n\theta\}}{4}, \frac{\{k_n\theta\}}{4}]}) \\ &= 2 \sin^2 \left( \frac{\{k_n\theta\}}{8} \right) + \frac{2}{3} \sin^4 \left( \frac{\{k_n\theta\}}{8} \right) + O \left( \sin^5 \left( \frac{\{k_n\theta\}}{8} \right) \right) \\ &= \frac{\{k_n\theta\}^2}{2^5} + \frac{\{k_n\theta\}^4}{2^{11} \cdot 3} + O(\{k_n\theta\}^5), \end{aligned} \tag{7}$$

since the mean curvature of the circle viewed inside  $\mathbb{R}^2$  is 1.

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

In view of the above equations, we see that

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

In view of the above equations, we see that  
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# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

In view of the above equations, we see that  
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# Exit time asymptotics for non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

In view of the above equations, we see that  
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the 'extrinsic diimension'  $d = 5$ ,  
and the 'mean curvature' is  $\frac{1}{2\sqrt{2}}$ .

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B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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All these give a good justification for developing a general theory of quantum stochastic geometry.

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

THANK

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

THANK  
YOU!!!

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Let  $\mathfrak{X} = \{A \in W^*(\mathbb{T}_\theta^2) \mid A = f_{-1}(U)V^{-1} + f_0(U) + f_1(U)V, f_1, f_0 \in L^\infty(\mathbb{T}), f_{-1}(t) := f_1(t + \theta)\}$ .

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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## Lemma

*The subspace  $\mathfrak{X}$  is closed in the ultraweak topology.*

# Exit time asymptotics on the non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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## Lemma

*The subspace  $\mathfrak{X}$  is closed in the ultraweak topology.*

## Proof.

Let  $A_\beta := f_{-1}^{(\beta)}(U)V^{-1} + f_0^{(\beta)}(U) + f_1^{(\beta)}(U)V$  be a convergent net in the ultraweak topology. Now  $\phi_1(A_\beta) = f_0^{(\beta)}(U)$ ,  $\phi_1(A_\beta V) = f_{-1}^{(\beta)}(U)$  and  $\phi_1(A_\beta V^{-1}) = f_1^{(\beta)}(U)$ . Since  $\phi_1$  is a normal map, which implies that  $f_0^{(\beta)}(U)$ ,  $f_1^{(\beta)}(U)$  and  $f_{-1}^{(\beta)}(U)$  (all of which are elements of  $L^\infty(\mathbb{T})$ ) are ultraweakly convergent, to  $f_0(U)$ ,  $f_1(U)$ ,  $f_{-1}(U)$  (say), and clearly  $f_{-1}(t) = \overline{f_1(t + \theta)}$ . □

# Exit time asymptotics on non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

## Lemma

Suppose  $f_1, f_0$  are as defined before and  $A \in \mathfrak{X}$ . Define

$$A_{s,t} := f_{-1}(e^{2\pi is} U) V^{-1} e^{-2\pi it} + f_0(e^{2\pi is} U) + f_1(e^{2\pi is} U) V e^{2\pi it}.$$

Suppose  $s, s' \in [0, 1)$  be such that  $|s - s'| \leq \frac{\epsilon}{4}$  where  $0 < \epsilon < \theta$ , and  $|\text{supp}(f_1)| < \epsilon$ , where  $|C|$  denotes the Lebesgue measure of a Borel subset  $C \subseteq \mathbb{R}$ . Then  $A_{s,t} \cdot A_{s',t'} \in \mathfrak{X}$ .

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## Proof.

It suffices to show that the coefficient of  $V^2$  in  $A_{s,t} \cdot A_{s',t'}$  is zero. By a direct computation, the coefficient of  $V^2$  is  $g(l) := f_1(s+l)f_1(s'+l-\theta)e^{2\pi i(t+t')}$ . But  $|(s+l) - (s'+l-\theta)| = |\theta + s - s'| > \epsilon$ . Now by hypothesis, we have  $|\text{supp}(f_1)| < \epsilon$ , so that  $f_1(s+l) \cdot f_1(s'+l-\theta) = 0$  and hence the lemma is proved.  $\square$

# Exit time asymptotics on non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

## Lemma

*Suppose  $A = f_{-1}(U)V^{-1} + f_0(U) + f_1(U)V$  and  $f_1(l)f_1(l + \theta) = 0$ , for  $l \in [0, 1)$ . Then  $A^{2n} \in \mathfrak{X}$ , for  $n \in \mathbb{N}$ .*

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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## Proof.

The coefficient of  $V^2$  in  $A^2$  is  $f_1(l)f_1(l + \theta)$  for  $l \in [0, 1)$  and this is zero by the hypothesis. Hence  $A^2 \in \mathfrak{X}$ . The coefficient of  $V$  in  $A^2$  is  $f_1^{(2)}(l) := f_1(f_0 + \tau_\theta(f_0))$ , where  $\tau_\theta$  is left translation by  $\theta$ . We have  $f_1^{(2)}(l)f_1^{(2)}(l + \theta) = 0$ , so that applying the same argument as before, we conclude that  $A^4 \in \mathfrak{X}$ . Proceeding like this we get the required result.  $\square$

# Exit time asymptotics on non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

Using the above three lemmas and von-Neumann's formula for minimum of two projections, we have

# Exit time asymptotics on non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
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Interplay  
between  
Geometry  
and  
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Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
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# Exit time asymptotics on non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
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# Exit time asymptotics on non-commutative 2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
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# Exit time asymptotics for the non-commutative 2-torus.

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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For  $l \in [0, 1)$ ,

B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
asymptotics  
on the non-  
commutative  
2-torus

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B.Das

Interplay  
between  
Geometry  
and  
Probability:

Exit time  
asymptotics  
of  
Brownian  
motion on  
manifolds.

Formulation  
of quantum  
exit time.

A case  
study: Exit  
time  
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on the non-  
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Geometry and  
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Exit time asymptotics of  
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Formulation  
of quantum  
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A case  
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Geometry  
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Exit time  
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Brownian  
motion on  
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A case  
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study: Exit  
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commutative  
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*For two projections  $A$  and  $B$  such that*

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B.Das

Interplay  
between  
Geometry  
and  
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Exit time  
asymptotics  
of  
Brownian  
motion on  
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Formulation  
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exit time.

A case  
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## Lemma

Let  $P = f_{-1}(U)V^{-1} + f_0(U) + f_1(U)V$  such that  $P$  is a projection and suppose  $f_0(t) = 0$  for some  $t$ . Then  $f_1(t) = f_1(t + \theta) = 0$ .